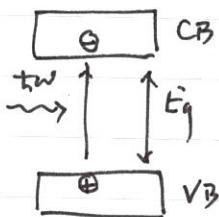


## Ch.14 Photodiode Detectors

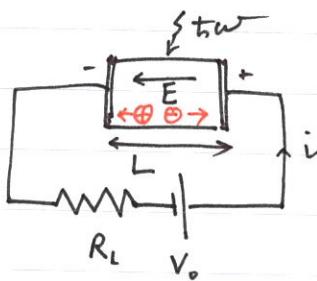
### - Problems in Photoconductive Detectors

1. Poor time response
2. Shot noise due to high dark current

### A. photoconductive Detectors (Ch.13)



• Absorption of light  $t_{\text{ow}}$   
 $\rightarrow e^-$  in CB : photoconductivity  
 $\rightarrow$  photocell

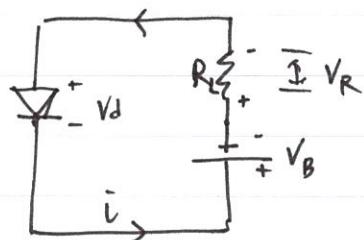


- $E = \frac{V_o}{L}$
- Due e-h pair generation,  
 $i_s$  creates :  $i_s = \frac{V_o}{R_L}$
- $i_s$  drops  $V_o$
- Drawback : background current  $i_o$  even w/o tow.  
 due to thermally generated e-h pairs.
- Practically  $i_o > i_s$  (inside the cell)

\* Photoconductive mode : p-n junction

(not photoconductive detector )

### (i) Photoconductive mode



(Fig. 14-2)

• Loop law (Kirchhoff)

$$\rightarrow V_d + V_B + V_R = 0$$

$$V_R = IR_L$$

$$\therefore i = -\frac{1}{R_L} (V_d + V_B)$$

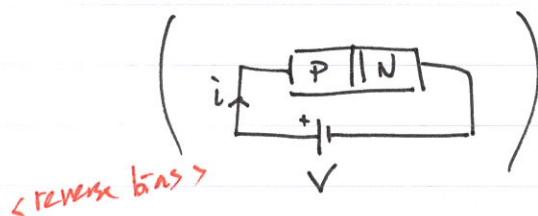
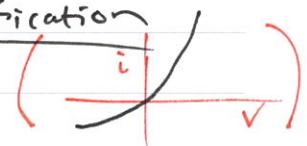
(load line)

• In p-n junction diode,

$$i = i_0 \left( e^{\frac{eV}{kT}} - 1 \right)$$

~~rectification~~

(14-21)



\* Why V applied : to balance the metal-semiconductor voltage drop in the circuit

• photocurrent : from e-h pairs generation via light absorption

$$i_\lambda = \frac{P_{in}}{hf} en$$

$$\rightarrow i = i_0 \left( e^{\frac{eV}{kT}} - 1 \right) - i_\lambda$$

shockley diode Eq.

i<sub>0</sub>: dark current

$$V_T = \frac{kT}{e}$$

(14-3)

• For no light

1. photovoltaic : i = 0

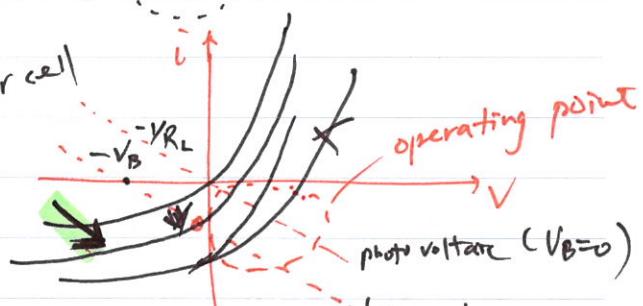
solar cell

2. photoconductive : i < 0

$$(i = -i_0)$$

dark current i<sub>0</sub>

Noisy



< Load line analysis >

### Application

1. photovoltaic  $\rightarrow$  solar cell

$$P_{\text{elec}} = i^2 R_L \quad , \quad V_d = -iR$$

for  $i_A \gg i_0$ ,

- $i \sim i_0 e^{(-e i R_L / \beta k_B T)} - i_A$  : solar cell current

### Efficiency

- $\eta_{sc} = \frac{P_{\text{elec}}}{P_{\text{in}}} = \frac{i^2 R_L}{P_{\text{in}}} \quad (\text{max must be optimized by } R_L)$

ex) Si solar cell area:  $4 \text{ cm}^2$

reverse saturation current density:  $1.5 \times 10^{-8} \text{ A/cm}^2$

diode ideal factor  $B$ :  $B=1$

light intensity:  $1000 \text{ W/m}^2$

$\lambda = 500 \text{ nm}$

$\eta_{\text{abs}} = 80\%$

in a solar cell

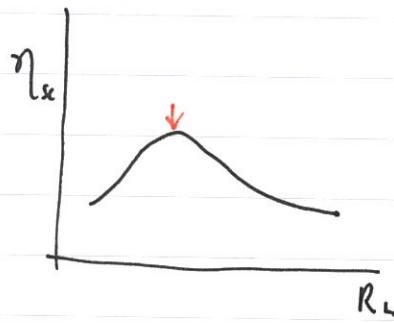
Q. What is optimum  $R_L$  & power conversion efficiency?

$$\text{Sol) } i_A = \frac{P_{\text{in}}}{h f} = \frac{P_{\text{in}} \lambda}{h c} e^{\eta} \Big|_{\text{ab}} = \frac{(0.4)(1000)(5 \cdot 10^{-7})(1.6 \cdot 10^{-19})(0.8)}{(6.63 \times 10^{-34})(3 \cdot 10^8)} = 0.129 \text{ A}$$

$$V_T = \frac{k_B T}{e} = \frac{(1.38 \times 10^{-23})(293)}{1.6 \times 10^{-19}} = 0.0253 \text{ V} \quad \text{at } T = 20^\circ\text{C} = 293\text{K}$$

$$\therefore i = i_0 e^{(-e i R_L / \beta k_B T)} - i_A = (6 \times 10^{-8}) e^{-\frac{i R_L}{0.0253}} - 0.129$$

\*  $\rightarrow$  need numerical calculation to find  $R_L$  for maximum efficiency ( $\eta_{sc} = i^2 R_L / P_{\text{in}}$ )



## 14-2 Output Saturation

### (ii) Photo voltaic mode

$$i = i_o \left[ e^{-\frac{V_d}{\beta V_T}} - 1 \right] - i_\lambda, \quad V_T = \frac{k_B T}{e}$$

for  $R_L \gg 1$ ,  $i \rightarrow 0$  (from  $i-V$  curve)

$$\therefore V_d = \beta V_T \ln \left( 1 + \frac{i_\lambda}{i_o} \right), \quad i_\lambda = \frac{P_{in} e \eta_{ab}}{h f}$$

① If  $i_\lambda \gg i_o$ ,

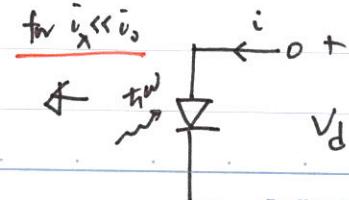
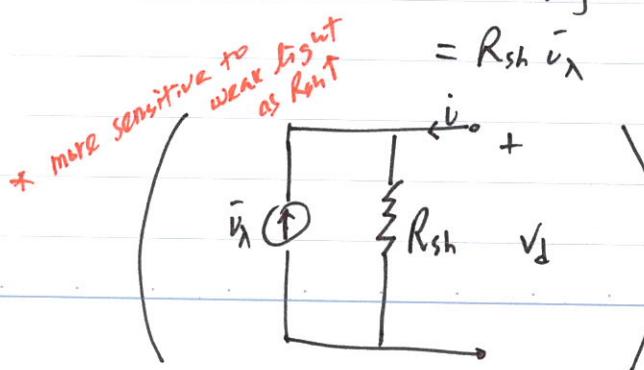
$$V_d \sim \beta V_T \ln \left( \frac{P_{in} e \eta_{ab}}{i_o h f} \right)$$

② If  $i_\lambda \ll i_o$ ,

$$V_d \sim \beta V_T \frac{P_{in} e \eta_{ab}}{i_o h f}$$

Taylor Expansion:  
 $\ln(1+x) = x - \frac{x^2}{2} + \frac{x^3}{3} - \dots$

$$= R_{sh} i_\lambda, \quad R_{sh} = \frac{\beta V_T}{i_o} \quad (\text{shunt resistance})$$



## 14.3 Response time

$$V_{out} = V_{max} (1 - e^{-t/RC}) , \quad RC = \tau : \text{time const.}$$

- Rise time  $t_r$ , ( $t_r = 10\% \rightarrow 90\% \text{ for } V_{out}$ )

$$t_r = (\ln 9) RC = 2.2 RC$$

→ determines the response time to modulation.

- Bandwidth  $B$ , (3-dB electrical bandwidth)

$$B = \frac{1}{2\pi t_r C}$$

$$= \frac{2}{2\pi t_r} = \frac{0.35}{t_r} \quad (3\text{-dB bandwidth})$$

Analysis:

The smaller  $C$ , the larger  $B$ .  
To reduce  $C$ ,

Techniques: ① Reverse bias  $V_B$  increase

→ photoconductive mode ( $V_B > 0$ )

② decrease the junction area

→ smaller detector area

→ Sensitivity vs. Speed trade off

③ Density of donor decreases

④ lower  $R_L$

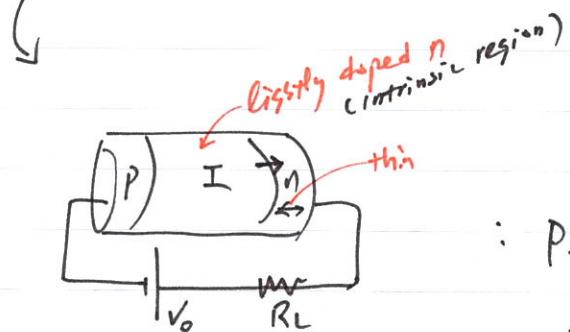
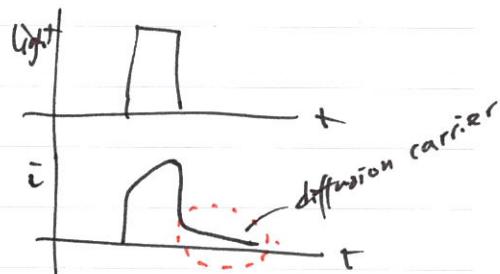
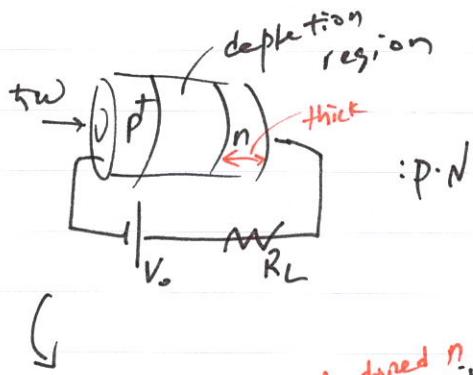
→ lower detector sensitivity, too.

## 14-4 Types of Photodiodes (PD)

### 1. PIN PD

- To solve diffusion carriers

→ donor concentration decrease in n-region



— most common today

(drift ~ diffusion)

ex) Si-PIN, Intrinsic region thickness: 0.1 mm

What is minimum  $t_r$ ,  $B$ , ~~reverse Bias~~?

Sol)  $v_s = 10^5 \text{ m/s}$ ,

$$t_r = \frac{d}{v_s} = \frac{10^{-4}}{10^5} = 10^{-9} \text{ s} = 1 \text{ ns}$$

$$B = \frac{0.35}{t_r} = \frac{0.35}{10^{-9}} = 350 \text{ MHz}$$

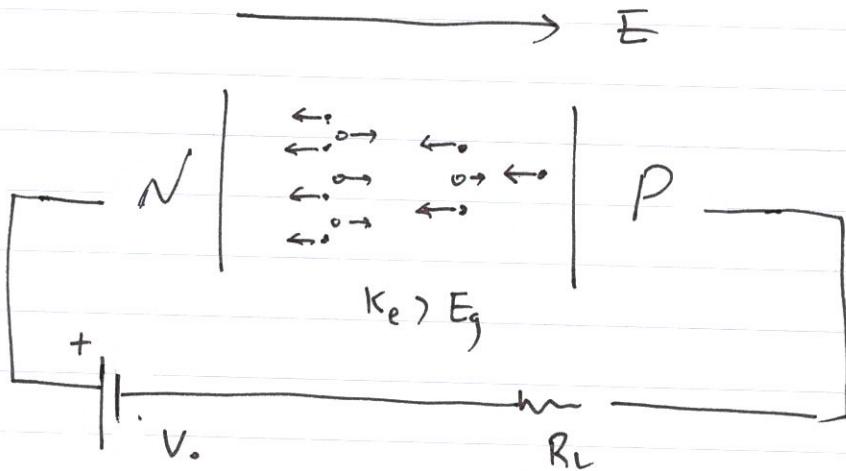
## 2. Avalanche PD

- PIN : small  $R_L$  for high speed  
 $\rightarrow$  small signal

Sel) multiplication in the depletion region

- analog of  $e^-$  multiplication in PMT.

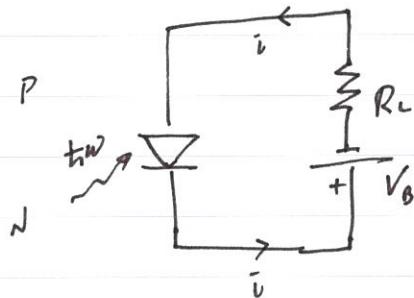
$\rightarrow$  In a P-N junction applied by a reverse bias  $V_r$ ,  
 $e^-$ s whose kinetic energy is enough generate  
additional  $e^-$ s via impact ionization process.  
 $\rightarrow$  avalanche photomultiplication



① For  $\Delta x$  moving,  $e^-$  lose potential energy  
by  $\Delta \times E \cdot e$ , but gains  
kinetic energy  $K$  ( $K > E_g$ )

② If  $K > E_g$ ,  $e^-$  creates e-h pair by collision

### 14-5 Signal to Noise ratio



(Reverse p-n)

$$P_{\text{sig}} = i_{\lambda}^2 R_L = \left( \frac{P_{\text{in}} \eta_{ab} e}{h f} \right)^2 R_L$$

$$i_{\lambda} = \frac{P_{\text{in}}}{h f} e \eta_{ab} \quad (14-2)$$

$\times$  Signal current ( $i_{\lambda}$ )

: measured current - dark current  
 $i_0$

(i) Shot noise induced electrical noise power :

$$P_{\text{shot}} = (i_N)^2 R_L = 2e(i_{\lambda} + i_0)BR_L$$

$$i_N = \sqrt{2eBR} \quad (13-29)$$

(ii) Thermal noise induced electrical noise power :

$$P_{\text{therm}} = \frac{V_N^2}{R_L} = 4k_B T B$$

$$V_N = \sqrt{4k_B T R_L B} \quad (13-33)$$

Due to  $R_L$  dependence on the shot noise,

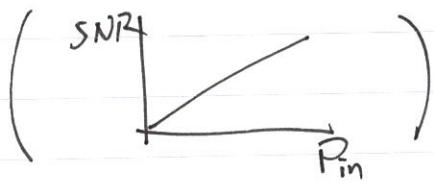
- $P_{\text{shot}}$  becomes a dominant noise source for large  $R_L$ ,  
for small  $R_L$ .
- $P_{\text{therm}}$  "

$$\text{SNR} \equiv \frac{P_{\text{sig}}}{P_{\text{shot}} + P_{\text{therm}}} = \frac{i_{\lambda}^2 R_L}{2e(i_{\lambda} + i_0)BR_L + 4k_B T B}$$

$\sqrt{\text{SNR}} \Rightarrow$  Amplitude  $\propto$  version.

(i) Large Sig:  $\bar{i}_\lambda \gg \bar{i}_o$  &  $\bar{i}_\lambda R_L \gg V_T$ ,  $V_T = \frac{k_B T}{e}$

$$\cdot SNR \sim \frac{\bar{i}_\lambda^2 R_L}{2e\bar{i}_o B R_L} = \frac{\bar{i}_\lambda}{2eB} \quad (\text{shot noise dominant})$$



(ii) Small Sig, small  $R_L$ :  $\bar{i}_\lambda \ll \bar{i}_o$  &  $\bar{i}_o R_L \ll V_T$

$$\cdot SNR \sim \frac{\bar{i}_\lambda^2 R_L}{4k_B T B}$$

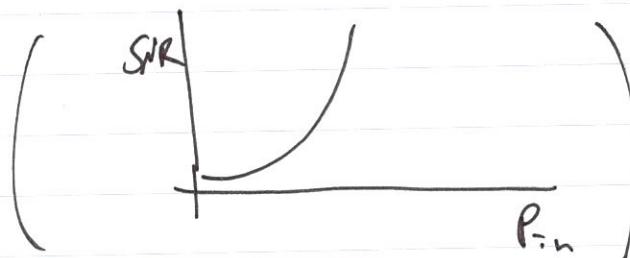
→ trade-off b/w response time and SNR

As  $R_L \uparrow$ , response time  $\uparrow$  (degrades in B)  
(due to RC)

(iii) Small Sig, large  $R_L$ :  $\bar{i}_\lambda \ll \bar{i}_o$  &  $\bar{i}_o R_L \gg V_T$

$$\cdot SNR \sim \frac{\bar{i}_\lambda^2 R_L}{2e\bar{i}_o B R_L} = \frac{\bar{i}_\lambda^2}{2e\bar{i}_o B}$$

→ limited by shot noise from dark current  $\bar{i}_o$



• Noise equivalent power :  $SNR = 1$   
(NEP)

• dB<sub>m</sub> :  $dB \rightarrow 1 \text{ mW}$  → denote optical power For power 1mW

Optical power in dBm =  $10 \log_{10} \left( \frac{P}{1 \text{ mW}} \right)$

ex)  $P = -20 \text{ dBm}$

$$\rightarrow -20 = 10 \log_{10} \left( \frac{P}{\text{mW}} \right)$$

$$P(\text{mW}) = 10^{-2} = 0.01 \text{ (mW)}$$

$P = +20 \text{ dBm}$

$$\rightarrow 20 = 10 \log_{10} (P)$$

$$P(\text{mW}) = 10^2 = 100 \text{ (mW)}$$